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Using x and θ as the coordinates, the kinetic and potential energies of the system can be expressed as

$$T = \frac{1}{2} m \dot{x}^2 + \frac{1}{2} J_0 \dot{\theta}^2 \quad (1)$$

where $J_0 = m (\ell + x)^2$ and

$$V = \frac{1}{2} k (x + \delta_{st})^2 - m g (\ell + x) \cos \theta \quad (2)$$

where $\delta_{st} = \frac{mg}{k}$. Equations (1) and (2) give

$$\frac{\partial T}{\partial \dot{x}} = m \dot{x} \quad ; \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{x}} \right) = m \ddot{x}$$

$$\frac{\partial T}{\partial \dot{\theta}} = J_0 \dot{\theta} \quad ; \quad \frac{d}{dt} \left(\frac{\partial T}{\partial \dot{\theta}} \right) = \dot{J}_0 \dot{\theta} + J_0 \ddot{\theta} = 2 m (\ell + x) \dot{x} \dot{\theta} + J_0 \ddot{\theta}$$

$$\frac{\partial T}{\partial x} = m (\ell + x) \dot{\theta}^2 \quad ; \quad \frac{\partial T}{\partial \theta} = 0$$

$$\frac{\partial V}{\partial x} = k (x + \delta_{st}) - m g \cos \theta \quad ; \quad \frac{\partial V}{\partial \theta} = m g (\ell + x) \sin \theta$$

The equations of motion can be derived using Lagrange's equations, Eq. (6.44), as:

$$m \ddot{x} - m (\ell + x) \dot{\theta}^2 + k x + m g - m g \cos \theta = 0 \quad (3)$$

$$m (\ell + x)^2 \ddot{\theta} + 2 m (\ell + x) \dot{x} \dot{\theta} + m g (\ell + x) \sin \theta = 0 \quad (4)$$

Using $\sin \theta \approx \theta$, $\cos \theta \approx 1$, and neglecting nonlinear terms involving $x^2 \ddot{\theta}$, $\dot{\theta}^2$, $x \theta$, and $\dot{x} \dot{\theta}$, Eqs. (3) and (4) can be reduced (linearized) to obtain:

$$m \ddot{x} + k x = 0 \quad (5)$$

$$m \ell^2 \ddot{\theta} + m g \ell \theta = 0 \quad (6)$$

Equations (5) and (6) correspond to the natural frequencies:

$$\omega_{n1} = \sqrt{\frac{k}{m}} \quad (7)$$

$$\omega_{n2} = \sqrt{\frac{g}{\ell}} \quad (8)$$

